

Quantization

classical mechanics (Poisson algebra, $\cdot, \{ \}$)

quantum mechanics (operator algebra, $*_{\hbar}$)

satisfying correspondence principle

$$\frac{1}{\sqrt{-1}\hbar} [\cdot]_{*_{\hbar}} \xrightarrow{(\hbar \rightarrow 0)} \{ \cdot \}$$

Poisson alg. $\xrightarrow{\quad}$ Operator alg. (non-comm. assoc. alg.)

$$f \xrightarrow{\quad} P_f^{(\hbar)}$$

$$\text{s.t. } \frac{1}{\sqrt{-1}\hbar} [P_f^{(\hbar)}, P_g^{(\hbar)}] \xrightarrow{(\hbar \rightarrow 0)} \{ f, g \}$$

* non-comm. assoc. alg. $\xrightarrow{\hbar \rightarrow 0}$ comm. assoc. alg.

↓
Space = spec

Kohn-Nirenberg quantization

$$(Q, P) = (q_1, \dots, q_n, p_1, \dots, p_n)$$

... Darboux coord.

$$\{q_i, p_j\} = \delta_{ij} \mapsto \frac{1}{i\hbar} [\hat{q}_i, \hat{p}_j] = \delta_{ij} \quad (H)$$

$\hat{q}_i: f(Q) \mapsto q_i \cdot f(Q)$... multiplication op.

$\hat{p}_j: f(Q) \mapsto \frac{\hbar}{i} \partial_{q_j} f(Q)$... differential op.

REM. Problem of domains of these operators is delicate.

$$\begin{aligned} \iota: \text{"Func"} &\longrightarrow A := \langle \hat{Q}, \hat{P} \rangle / (H) \\ \varphi &\longrightarrow \iota(\varphi) \end{aligned}$$

$$\varphi * \psi := \iota^{-1}(\iota(\varphi) \circ \iota(\psi))$$

Roughly speaking, this procedure gives non-commutative assoc. product $*$ on "Func."

Precisely, appropriate symbol class

$$\mathcal{L}_{KN} : \text{Func.} \xrightarrow{\mathcal{F}} A := \Psi DO.$$

$$p^\alpha q^\beta \longmapsto \hat{q}^\beta \hat{p}^\alpha$$

$$\frac{\hbar}{\sqrt{-1}} \partial_q f(q) \xrightleftharpoons[\mathcal{F}^{-1}]{\mathcal{F}} p \cdot \hat{f}(p)$$

Fourier transf.

$$\mathcal{L}_{KN} (\varphi(p, q)) f(q)$$

$$:= \int \varphi(q, p) e^{\frac{\sqrt{-1}}{\hbar} (q - q') \cdot p} f(q') \, d'q' \, dp$$

$$(d := \frac{1}{\sqrt{2\pi\hbar}^n} d)$$

$P_\varphi := \mathcal{L}_{KN}(\varphi(p, q)) \dots$ pseudo-differential op. with symbol φ .

TH (symbol calculus)

$$P_\varphi \circ P_\psi = P(\varphi *_{KN} \psi),$$

where

$$\varphi *_{KN} \psi = \sum_{\alpha: \text{multi index}} \frac{(-\sqrt{-1}\hbar)^{|\alpha|}}{\alpha!} \partial_p^\alpha \varphi(p, q) \cdot \partial_q^\alpha \psi(p, q).$$

Furthermore, we have

$$\text{symbol of } \frac{1}{\sqrt{-1}\hbar} [P_\varphi, P_\psi] \xrightarrow{(\hbar \rightarrow 0)} \{\varphi, \psi\}.$$

Proof. $((P\varphi \circ P\psi) f) (q)$

$$= \int \varphi(q, p) e^{\frac{i}{\hbar}(q-q') \cdot p} \times \left(\int \psi(q', p') e^{\frac{i}{\hbar}(q'-q'') p'} f(q'') dq'' dp' \right) dq dp$$

$$= \iint \varphi(q, p) \psi(q', p') \times e^{\frac{i}{\hbar} [(q-q'') p' + (q'-q)(p'-p)]} dq' dp \times f(q'') dq'' dp'$$

Taylor
$$= \iint \left(\sum_{\alpha} \frac{(p-p')^{\alpha}}{\alpha!} \partial_{p'}^{\alpha} \varphi(q, p') \right) \times \psi(q', p') e^{\frac{i}{\hbar}(q'-q)(p'-p)} dq dp \times e^{\frac{i}{\hbar}(q-q'') p'} f(q'') dq'' dp'$$

$$= \iint \left(\sum_{\alpha} \frac{1}{\alpha!} \partial_{p'}^{\alpha} \varphi(q, p') \psi(q', p') \cdot \left(-\frac{\hbar}{i} \partial_{q'} \right)^{\alpha} e^{\frac{i}{\hbar}(q'-q)(p'-p)} \right) dq dp \times e^{\frac{i}{\hbar}(q-q'') p'} f(q'') dq'' dp'$$

integral by parts
$$= \iint \left(\sum_{\alpha} \frac{(\frac{\hbar}{i})^{|\alpha|}}{\alpha!} \partial_{p'}^{\alpha} \varphi(q, p') \cdot \partial_{q'}^{\alpha} \psi(q', p') \cdot e^{\frac{i}{\hbar}(q'-q)(p'-p)} \right) dq dp \times e^{\frac{i}{\hbar}(q-q'') p'} f(q'') dq'' dp'$$

Fourier inversion formula
$$= \left(\sum_{\alpha} \frac{(-i\hbar)^{|\alpha|}}{\alpha!} \partial_{p'}^{\alpha} \varphi(q, p') \partial_{q'}^{\alpha} \psi(q, p') \right) \times e^{\frac{i}{\hbar}(q-q'') p'} f(q'') dq'' dp'$$

$$= (P\varphi *_{KN} \psi f) (q)$$

REM. \int should be replaced with $os-\int$
 Thus, the calculus above is formal.
 In general, $P\varphi$ is not Hermitian.

Weyl quantization

$$z_w : \text{"Func"} \longrightarrow A$$

$$\varphi \longmapsto W\varphi$$

where

$$W\varphi := \int \hat{\varphi}(\eta, \xi) e^{i(\eta \hat{q} + \xi \hat{p})} d\eta d\xi$$

... Weyl type ψ b.o.

Then we have

TH (Weyl calculus, Moyal product)

$$W\varphi \circ W\psi = W\varphi *_{\hbar} \psi,$$

where

$$\varphi *_{\hbar} \psi = \sum_{\alpha, \beta} \frac{\left(\frac{i\hbar}{2}\right)^{|\alpha+\beta|}}{\alpha! \beta!} \partial_p^\alpha \partial_q^\beta \varphi(p, q) \cdot \partial_q^\alpha (\partial_p)^\beta \psi(q, p)$$

Furthermore,

$$\text{Symbol of } \frac{1}{\hbar \sqrt{-1}} [W\varphi, W\psi] \xrightarrow{\hbar \rightarrow 0} \{\varphi, \psi\}.$$

Proof

$$\mathcal{F}: \mathbb{R}^{2n}_{(q,p)} \longrightarrow \hat{\varphi} \in \mathbb{R}^{2n}_{(\eta,\xi)} ; \text{Fourier transf.}$$

$$P := \hat{p}, \quad Q := \hat{q}$$

$$\iint \hat{\varphi}(\eta_1, \xi_1) \hat{\psi}(\eta_2, \xi_2) e^{i(\eta_1 Q + \xi_1 P)} \cdot e^{i(\eta_2 Q + \xi_2 P)} d\eta_1 d\xi_1 d\eta_2 d\xi_2$$

$$\stackrel{\text{BCH}}{=} \iint \hat{\varphi}(\eta_1, \xi_1) \hat{\psi}(\eta_2, \xi_2) \cdot e^{\frac{i\hbar}{2}(\xi_1 \eta_2 - \xi_2 \eta_1)} \cdot e^{i(\xi P + \eta Q)} d\eta_1 d\xi_1 d\eta_2 d\xi_2$$

can be seen as
group 2 cocycle

$\begin{cases} \xi = \xi_1 + \xi_2 \\ \eta = \eta_1 + \eta_2 \end{cases}$

$$\stackrel{\text{Taylor}}{=} \iint \hat{\varphi}(\eta_1, \xi_1) \hat{\psi}(\eta_2, \xi_2) \left(\sum_n \frac{1}{n!} \left(\frac{i\hbar}{2} \right)^n (\xi_1 \eta_2 - \xi_2 \eta_1)^n \right) \times e^{i(\xi P + \eta Q)} d\eta_1 d\xi_1 d\eta_2 d\xi_2$$

$$= \iint \sum_{\alpha, \beta} \frac{\left(\frac{i\hbar}{2} \right)^{|\alpha+\beta|}}{\alpha! \beta!} \xi_1^\alpha \eta_1^\beta \hat{\varphi}(\eta_1, \xi_1) \cdot \eta_2^\alpha (-\xi_2)^\beta \hat{\psi}(\eta_2, \xi_2) \times e^{i(\xi P + \eta Q)} d\eta_1 d\xi_1 d\eta_2 d\xi_2$$

$$= \iint \sum_{\alpha, \beta} \frac{\left(\frac{i\hbar}{2} \right)^{|\alpha+\beta|}}{\alpha! \beta!} \overbrace{(-i\partial_p)^\alpha (-i\partial_q)^\beta \varphi(\eta, \xi)} \cdot \overbrace{(-i\partial_q)^\alpha (i\partial_p)^\beta \psi(\eta-\eta_1, \xi-\xi_1)} \times e^{i(\xi P + \eta Q)} d\eta_1 d\xi_1 d\eta d\xi$$

convolution

$$= \int \sum_{\alpha, \beta} \frac{\left(\frac{i\hbar}{2} \right)^{|\alpha+\beta|}}{\alpha! \beta!} \overbrace{(\partial_q^\beta \partial_p^\alpha \varphi(q,p) \cdot \partial_p^\beta (-\partial_q)^\alpha \psi(q,p))} (\eta, \xi) \times e^{i(\xi P + \eta Q)} d\eta d\xi$$

$$= \int \widehat{\varphi *_{\hbar} \psi} (\eta, \xi) e^{i(\xi P + \eta Q)} d\eta d\xi$$

$$\therefore W\varphi \circ W\psi = W\varphi *_{\hbar} \psi.$$

Consider "formal" Weyl algebra bundle W_U on a local Darboux coordinates.

Glueing these algebra bundles, we would like to construct globally defined non-commutative associative algebra bundle which can be regarded as deformation quantization in some sense.

Deformation quantization

- Background information
of Quantization

(1) Existence (Deformation quantization)

$\forall (M, \omega)$: Symplectic manifold

$$\exists A := (C^\infty(M) \llbracket \hbar \rrbracket, *)$$

... Deformation quantization of
(M, ω)

- DeWilde - Lecomte
- Omori - Maeda - Yoshioka
- Fedosov
- Kontsevich

(2) classification

$$\{D, Q, \text{ on } (M, \omega)\} / \sim$$

$$\downarrow \\ \check{H}^2(M, \mathbb{R} \llbracket \hbar \rrbracket)$$

- Deligne
- Omori - Maeda - Miyazaki - Yoshioka
- Nest - Tsygán
- Kontsevich

⊙ Deformation quantization (def.)

DEF. (Bayen - Flato - Fronsdal

- Lichnerowicz - Sternheimer)

A deformation quantization/star product of a Poisson manifold $(M, \{, \})$ is a product $*$ on $C^\infty(M)[[v]]$, the space of formal power series of parameter v with coefficients in $C^\infty(M)$:

$$f * g = f \cdot g + v \pi_1(f, g) + \dots + v^n \pi_n(f, g) + \dots$$

satisfying

(1) $*$ is associative,

$$(2) \pi_1(f, g) = \frac{1}{2\sqrt{-1}} \{f, g\}$$

(3) each π_n is an $\mathbb{C}[[v]]$ -bilinear, bidifferential operator.

② Weyl manifold, contact Weyl manifold
for deformation quantization

DEF. W : Weyl algebra

$\overset{\text{def}}{\longleftrightarrow}$

binary operator: $*$

generators:

$$\nu, z_1 = x_1, \dots, z_n = x_n, z_{n+1} = y_1, \dots, z_{2n} = y_n$$

relations:

$$[x_i, x_j] = [y_i, y_j] = [\nu, x_i] = [\nu, y_j] = 0$$

$$[x_i, y_j] = -\nu \delta_{ij} \quad ([z_i, z_j] = \nu \Lambda_{ij}, \Lambda = \begin{bmatrix} 0 & I_n \\ -I_n & 0 \end{bmatrix})$$

where $[a, b] := a * b - b * a$.

$$\exists \bar{\cdot} : a \mapsto \bar{a}; \quad \overline{a * b} = \bar{b} * \bar{a}, \quad \bar{\nu} = -\nu, \quad \bar{z}_i = z_i$$

$$\exists d := \text{deg} : \nu^k z^\alpha \mapsto 2k + |\alpha|, \quad \alpha: \text{multi-index} \quad \square$$

DEF. C : contact Weyl algebra

$\overset{\text{def}}{\longleftrightarrow}$

$$C = \tau C + W$$

where W is a Weyl algebra with

$$[\tau, \nu] = 2\nu^2, \quad [\tau, z_i] = \nu z_i. \quad \square$$

② automorphism of algebra

DEF. Φ ; ν -automorphism of Weyl algebra W

\Leftrightarrow (1) Φ : \mathbb{C} -linear

(2) $\Phi(\nu) = \nu$

(3) $\Phi(a * b) = \Phi(a) * \Phi(b)$

(4) $\Phi(\bar{a}) = \overline{\Phi(a)}$ □

We consider the following algebra bundles:

$$W_M := \bigsqcup_{\nu} W_{\nu} / \sim \quad (W_{\nu} := U \times W)$$

$$C_M := \bigsqcup_{\nu} C_{\nu} / \sim \quad (C_{\nu} := U \times C)$$

where " \sim " is a suitable equivalence relation discussed below:

DEF. Let $U(x, y)$ be a Darboux chart of symplectic manifold M .

$$\begin{aligned} \# : C^{\infty}(U) [U] &\longrightarrow \Gamma(W_{\nu}) \subset \Gamma(C_{\nu}) \\ f &\longmapsto f^{\#} = \sum_{\alpha, \beta} \frac{x^{\alpha} y^{\beta}}{\alpha! \beta!} \partial_x^{\alpha} \partial_y^{\beta} f(x, y) \\ &\quad \dots (\text{local}) \text{ Weyl function} \end{aligned}$$

$$\mathcal{F}(W_{\nu}) := \{ f^{\#} \mid f \in C^{\infty}(U) [U] \} \quad \square$$

DEF. Φ : a "local" Weyl diffeo.

def \iff (1) $\Phi_z : W_z \longrightarrow W_{\phi(z)}$ is a ν -auto.
where $z \in M$ (base mfd.), ϕ is the local diffeo. induced by Φ .

$$(2) \Phi^*(\mathcal{F}(W_{\phi(v)})) = \mathcal{F}(W_v). \quad \square$$

PROP. Φ is a "local" Weyl diffeo.

$\implies \phi : U \rightarrow \phi(U)$ is a symplectic diffeo. \square

Conversely,

TH. ([OMY, 3.7]) For any symplectic diffeo. $\phi : U \rightarrow \phi(U)$, there exists a Weyl diffeo. Φ which induces ϕ . \square

This theorem plays a crucial role to glue trivial algebra bundle together preserving Weyl functions $\{\mathcal{F}(U)\}_U$. (i.e. $\perp W_U / \sim$)

DEF. $\Psi : C \rightarrow C$, ν -automorphism

$\stackrel{\text{def.}}{\iff}$ (1) Ψ : algebra isomorphism,

(2) $\Psi|_W$: ν -auto, of Weyl algebra. \square

DEF. $\Psi : C_\nu \rightarrow C_{\phi(\nu)}$

is a modified contact Weyl diffeo.

(MCWD for short.)

$\stackrel{\text{def.}}{\iff}$ (1) $\Psi_\mathbb{Z}$: ν -automorphism,

(2) $\Psi|_{W_\nu}$: Weyl diffeomorphism. \square

This notion is also important to glue trivial contact Weyl algebra bundle together.

⑩ Properties of MCWD

DEF. $\tilde{\tau}_v(z) \stackrel{\text{def}}{=} \tau + \sum_{i,j} z_i \omega_{ij} Z_j \in \Gamma(C_v)$

$$[\tilde{\tau}_v(z), F] = 2v^2 \partial_v F + v \sum_i (z_i)^\# \cdot \partial_{z_i} F$$

REM. $\text{ad}\left(\frac{1}{v}F\right)(a) := \begin{cases} \frac{1}{v}[F, a] & (a \in W) \\ 2F + \frac{1}{v}[F, \tau] & (a = \tau) \end{cases} \quad \square$

TH. $\forall \Psi$: modified contact Weyl diffeo.

$$\exists f \in C^\infty(V) \llbracket v \rrbracket, \exists a \in (v^2) \in C^\infty(V) \llbracket v \rrbracket$$

$$; \Psi^* \tilde{\tau}_{\phi(v)} = \tilde{\tau}_v + f^\# + \underline{a(v^2)} \quad \blacksquare$$

DEF. A modified contact Weyl diffeo. Ψ satisfying

$$\Psi^* \tilde{\tau}_{\phi(v)} = \tilde{\tau}_v + f^\#$$

is called a contact Weyl diffeomorphism.

(CWD, for short.) □

TH. [OMY, 4.7]

$\forall \phi$: local symplectic diffeomorphism

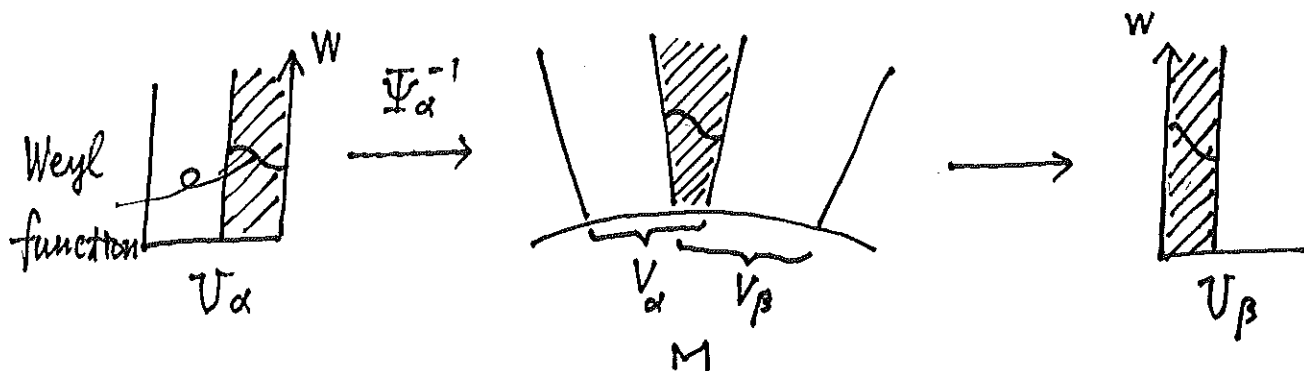
$\exists \hat{\phi} := \Psi$: contact Weyl diffeomorphism

; Ψ induces ϕ . □

DEF. (Weyl manifold)

$\{ W_M \xrightarrow{\pi} M = \bigcup_{\alpha} V_{\alpha}, \Psi_{\alpha}: \pi^{-1}(V_{\alpha}) \rightarrow W_{V_{\alpha}} := V_{\alpha} \times W \}$
 is called a Weyl manifold.

def. $\forall \Psi_{\alpha\beta} := \Psi_{\beta} \circ \Psi_{\alpha}^{-1}: W_{V_{\alpha}} \rightarrow W_{V_{\beta}}$ is a Weyl diffeo.



TH. [Omori - Maeda - Yoshioka]

$\forall (M, \omega)$: symplectic manifold

$\exists \{ W_M \xrightarrow{\pi} M \}$: Weyl manifold

; $M = \bigcup_{\alpha} V_{\alpha}, V_{\alpha}$: Darboux chart. ■

TH. [Omori - Maeda - Yoshioka - Miyazaki]

Assume that (M, ω) is a symplectic manifold.

Then,

$\{ \{ W_M \xrightarrow{\pi} M : \text{Weyl manifold} \} / \sim$

$\leftrightarrow \{ * : (\text{hermite type}) \text{ star product} \} / \sim$

$\leftrightarrow [\omega] + H^2(M, \nu^2 \mathbb{R} \llbracket \nu^2 \rrbracket)$ ■

● Automorphism.

We will treat

$$\text{Aut}(C_M) := \{ \Psi : C_M \rightarrow C_M \mid \text{global MCWD.} \}.$$

Before it. we prepare fundamental Proposition.

Prop. [Y. Prop. 2.5 and 2.18]

(1) $\Phi : W \rightarrow W$, ν -auto. of Weyl algebra

$$\Rightarrow \exists! A = (a_{ij}) \in Sp(n, \mathbb{R}), \exists! F \in W_{\frac{1}{3}}^{+,0}$$

$$\text{s.t. } \Phi = \hat{A} \circ e^{\frac{1}{\nu} \text{ad} F} \quad \left\{ a \in W \mid a = \sum_{2l+|\alpha| \geq 3} a_{\alpha} \nu^l z^{\alpha}, \bar{a} = a \right\}$$

$$\text{where } \hat{A} z^i = \sum a_{ij} z^j, \hat{A} \nu = \nu.$$

We call \hat{A} a Weyl lift of A .

(2) $\Psi : C \rightarrow C$, ν -auto. of contact Weyl algebra

$$\Rightarrow \exists! A = (a_{ij}) \in Sp(n, \mathbb{R}), \exists! F \in W_{\frac{1}{3}}^{+,0}, \exists! c(\nu^2) \in \mathbb{R} \llbracket \nu^2 \rrbracket$$

$$\text{s.t. } \Psi = \hat{A} \circ e^{\text{ad}(\frac{1}{\nu}(F + c(\nu^2)))}$$

Proof.

$$(1) \mathcal{M} := \left\{ a \in \sum_{2l+|\alpha| \geq 1} a_{l\alpha} v^l z^\alpha \in W \right\} \dots \text{unique maximal ideal.}$$

1st. $\Phi(\mathcal{M}) \subset \mathcal{M}$

$$\therefore \Phi(z^i) = \sum a_{ij} z^j + O(2) \dots (i)$$

where

$O(2)$ is the collection of the terms degree ≥ 2 .

Applying (i) to $[z^i, z^j] = v \Lambda^{ij}$, we see

$$A \in Sp(n, \mathbb{R}).$$

2nd $\hat{A}^{-1} \circ \Phi(z^i) = z^i + g_{(2)}^i + O(3) \dots (ii)$

where

$g_{(2)}^i$ is the collection of the terms of homogeneous degree = 2.

$O(3)$ is the collection of the terms degree ≥ 3 .

Applying (ii) to $[z^i, z^j] = v \Lambda^{ij}$, we have

$$v \partial_{i+n} g_{(2)}^j = v \Lambda^{il} \partial_l g_{(2)}^j = [z^i, g_{(2)}^j]$$

$$= [z^j, g_{(2)}^i] = v \Lambda^{jk} \partial_k g_{(2)}^i = v \partial_{j+n} g_{(2)}^i$$

This means "closedness", i.e.

$$d(g_{(2)}^j dz^{j+n} + g_{(2)}^i dz^{i+n}) = 0$$

Then, by Poincaré's Lemma,

$\exists 1 F_{(3)}$: homogeneous degree = 3

$$\text{s.t. } \frac{1}{\nu} [Z^i, F_{(3)}] = g_{(2)}^i$$

$$\therefore \hat{A}^{-1} \circ \Phi(Z^i) = e^{\frac{1}{\nu} F_{(3)}}(Z^i) + O(3).$$

3rd Repeating this process, we obtain

$$\Phi(Z^i) = \hat{A} \circ \underbrace{e^{\text{ad}(\frac{1}{\nu} F_{(3)})} \circ e^{\text{ad}(\frac{1}{\nu} F_{(4)})} \circ \dots \circ e^{\text{ad}(\frac{1}{\nu} F_{(k)})}}_{k \rightarrow \infty} \dots(Z^i)$$

$k \rightarrow \infty \downarrow$ Baker-Campbell-Hausdorff.

$$e^{\text{ad}(\frac{1}{\nu} F)} \quad (\exists F \in W_3^{0+})$$

$$\therefore \Phi = \hat{A} \circ e^{\text{ad}(\frac{1}{\nu} F)}$$

The uniqueness is inductively checked!

(2) By (1), we have $\Psi|_W = \hat{A} \circ e^{\text{ad}(\frac{1}{v}F)} =: \psi$ on C .
... (iii)

$$\therefore \psi^{-1} \circ \Psi(z^i) = z^i \dots \text{(iv)}$$

Applying (iv) to $[\tau, z^i] = v z^i$.

$$\psi^{-1} \circ \Psi(\tau) = \tau + b \quad (\exists b = b_0 + v^2 b_2 + \dots + v^{2k} b_{2k} + \dots \in \mathbb{R} \llbracket v^2 \rrbracket).$$

Put

$$c(v^2) := \sum v^{2k} \frac{b_{2k}}{2(1-2k)},$$

then we obtain

$$e^{\text{ad}(\frac{1}{v}c(v^2))} z^i = z^i.$$

$$\therefore e^{\text{ad}(-\frac{1}{v}c(v^2))} \circ \psi^{-1} \circ \Psi = 1 \text{ on } C.$$

$$\therefore \Psi \stackrel{\text{(iii)}}{=} A \circ e^{\text{ad}(\frac{1}{v}(F + c(v^2)))}.$$

Summarizing what mentioned above,

(i) Ψ : v -auto. of contact Weyl alg. C

$$\Rightarrow \Psi = \hat{A} \circ \exp\left(\frac{1}{v}(c(v^2) + F)\right)$$

$$(\exists A \in Sp(n, \mathbb{R}), \exists c(v^2) \in \mathbb{R} \llbracket v^2 \rrbracket, \exists F \in W)$$

(ii) Ψ : MCWD which induces the id. map on the base space

$$\Rightarrow \Psi = \mathbb{1} \circ \exp\left(\frac{1}{v}(g(v^2) + v^2 f^\#)\right)$$

$$(\exists g(v^2) \in C^\infty(M) \llbracket v^2 \rrbracket, f^\# \in \mathcal{F}_M)$$

(iii) Ψ : MCWD

$$\Rightarrow \Psi = \tilde{\psi} \circ \exp\left(\frac{1}{v}(g(v^2) + v^2 f^\#)\right)$$

proof of (iii)

$$\Psi \xrightarrow{\substack{\text{induce a symp.} \\ \text{diffeo. on } M.}} \psi_M \xrightarrow{\text{canonical MCWD lift}} \hat{\psi}_M$$

Then,

$$\hat{\psi}_M \circ \Psi \in \underline{Aut}(C_M)$$

$$\stackrel{(ii)}{\Rightarrow} \hat{\psi}_M^{-1} \circ \Psi = \exp\left(\frac{1}{v}(g(v^2) + v^2 f^\#)\right)$$

⑩ Automorphism (Group)

DEF. $\text{Aut}(C_M) := \left\{ \Psi: C_M \rightarrow C_M \mid \begin{array}{l} \text{fiberwise } \nu\text{-auto} \\ \Psi^*(F_M) = F_M \end{array} \right\}$

Aut(C_M) := $\left\{ \Psi \in \text{Aut}(C_M) \mid \begin{array}{l} \Psi \text{ induces the} \\ \text{identity map} \\ \text{on the base space} \end{array} \right\}$

We would like to consider the following sequence:

$$1 \rightarrow \underline{\text{Aut}}(C_M) \rightarrow \text{Aut}(C_M) \rightarrow \text{Diff}(M, \omega) \rightarrow 1.$$

Before studying the above sequence, we would like to note several basics relating to infinite dimensional manifold and infinite dimensional Lie group.

Fact (Omori - de la Haap)

G Banach-Lie group

effective

M : finite dimensional smooth manifold

$\Rightarrow \dim G < \infty$.

□

In order to treat $\text{Aut}(C_M)$ and $\text{Diff}(M, \omega)$,
as Lie groups, we need a framework
of infinite dimensional calculus,
manifolds, and Lie groups beyond
Banach-category!

↓

Difficulties (which we need to overcome)

(1) How can we construct infinite
dimensional calculus?

What is smoothness?

(2) Definitions of manifold,
(co) tangent space, Lie group,
exponential.

As to (1) (cf. [KM]... Kriegl-Michor)

- smoothness of curves \rightarrow no problem!

- smoothness of maps

"smooth curves" $\overset{\downarrow}{\mapsto}$ "smooth curves"

① convenient space E

$\stackrel{\text{def}}{\iff} \forall c_1 \in C^\infty(\mathbb{R}, E), \exists c_2 \in C^\infty(\mathbb{R}, E); c_2' = c_1$

$\iff \forall c: \mathbb{R} \rightarrow E (\forall \ell \in E^*; \ell \circ c: \text{smooth})$
 $\iff c: \text{smooth}$

$\iff \forall B: \text{bdd. closed convex set } \subset E$

; $E_B := \text{linear span of } B \text{ is Banach}$
with norm

$$p_B(v) := \inf \{ \lambda > 0; v \in \lambda B \}.$$

② Everything else is a theorem!

- linearity, - Leibniz rule

- chain rule, - mean value

- Taylor, - the fundamental theorem

- cartesian closedness

$$(\text{i.e. } C^\infty(U, C^\infty(V, E)) \cong C^\infty(U \times V, E).)$$

③ The definition of manifolds modeled on convenient spaces is given in the standard manner.

⑩ Tangent vectors / bundles

We investigate tangent vectors seen as kinematic tangent vector (via curves), and then "push forward" is defined in the usual manner.

⑪ Differential forms

There is only one suitable class satisfying all requirements (cf [KM, P351]).

$$\Omega^k(M) := C^\infty \left(\underbrace{L^k_{alt}}_{\text{alternative}}(TM, M \times \mathbb{R}) \right)$$

↓

$\Omega^k(M)$ is closed under pull-back f^* , wedge product \wedge , exterior derivative d , interior product ι_x , Lie derivative L_x !

(2) Lie groups

The definition of "Lie group" is given in the usual manner.

⊙ regular Lie groups

regularity introduced by



- Omoni-Maeda-Yoshioka-Kobayashi '82
- Milnor '84

"smooth curves in Lie algebras
integrate to
smooth curves in Lie groups
in a smooth way"

⊙ Precise definition

$$\forall X \in C^\infty(\mathbb{R}, \mathfrak{g}) \quad (\mathfrak{g} := \text{Lie } G)$$

$$\exists g \in C^\infty(\mathbb{R}, G)$$

$$\text{s.t. } \partial_t g(t) = R|_{g(t)} X(t).$$

REM. $X(t) := tX \implies g(t) = \exp(tX).$

Th. (i) $\text{Aut}(C_M)$ is a Lie group
modeled on

$$C := C^\infty(M) \llbracket V^2 \rrbracket + V^2 \mathcal{F}_M.$$

(ii) $\text{Aut}(C_M)$ is a Lie group
modeled on

$$\mathcal{X}(M, \omega) \times C, \text{ where}$$

$\mathcal{X}(M, \omega)$ is the Lie algebra of
all symplectic vector fields.

Combining the above theorem, with Th. 4.7. in [OMY], we obtain,

Prop. $\forall \Psi \in \text{Aut}(C_M)$

$\exists \tilde{\psi} : \text{MCWD}$ (canonical lift of ψ)

$\exists g(v^2) + v^2 f^\#(v^2)$

s.t. $\Psi = \tilde{\psi} \circ e^{\text{ad}(\frac{1}{v}(g(v^2) + v^2 f^\#(v^2)))}$.

Then we have

$\text{Aut}(C_M) \longleftrightarrow \text{Diff}(M, \omega) \times \mathbb{C}$

$\Psi \longleftrightarrow (\psi, g(v^2) + v^2 f^\#(v^2))$

Proof.

$\Psi \rightsquigarrow \psi : \text{induced symplectic diffeo.} \rightsquigarrow \tilde{\psi} : \text{lift}$

$\cdot \tilde{\psi}^{-1} \circ \Psi \in \text{Aut}(C_M)$.

$\therefore \tilde{\psi}^{-1} \circ \Psi \stackrel{\text{Th.}}{=} e^{\frac{1}{v} \text{ad}(g(v^2) + v^2 f^\#(v^2))}$ ■

Fundamental Theorem.

$$(1) \quad 1 \rightarrow N \rightarrow G \rightarrow M \rightarrow 1$$

is an exact sequence of Fréchet-Lie groups.

- $\exists M \supset U \xrightarrow{s} M$, local smooth section

- M, N : regular

$\Rightarrow G$: regular □

(2) Assume that M be compact. Then,

$$\text{Diff}(M, \omega) = \{ \varphi : \text{symplectic diffeo. on } M \}$$

is a regular Fréchet-Lie group

modeled on

$$\mathcal{X}(M, \omega) = \{ X : \text{symplectic vector field on } M. \} \quad \square$$

Combining the above Prop. and Fundamental Theorem, we have,

Th. (1) $1 \rightarrow \underline{\text{Aut}}(C_M) \rightarrow \text{Aut}(C_M) \rightarrow \text{Diff}(M, \omega) \rightarrow 1$

is an exact sequence of Fréchet-Lie groups, where (M, ω) is a compact symplectic manifold.

(2) $\underline{\text{Aut}}(C_M)$ is regular.

(3) $\text{Diff}(M, \omega)$ is regular.

(4) $\text{Aut}(C_M)$ is regular.

Proof.

We first show $\text{Aut}(C_M)$ is a Fréchet-Lie group modeled on $\mathcal{X}(M, \omega) \times C$.

Next, we consider composition.

$$\begin{aligned}\Psi_1 \circ \Psi_2 &= \tilde{\psi}_1 \circ e^{\text{ad}(\frac{1}{v}H_1)} \circ \tilde{\psi}_2 \circ e^{\text{ad}(\frac{1}{v}H_2)} \\ &= \widetilde{\psi_1 \circ \psi_2} \circ \left\{ (\widetilde{\psi_1 \circ \psi_2})^{-1} \circ \tilde{\psi}_1 \circ \tilde{\psi}_2 \right\} \\ &\quad \circ \left\{ \tilde{\psi}_2^{-1} \circ e^{\text{ad}(\frac{1}{v}H_1)} \circ \tilde{\psi}_2 \right\} \circ e^{\text{ad}(\frac{1}{v}H_2)} \\ &= \widetilde{\psi_1 \circ \psi_2} \circ e^{\text{ad}(\frac{1}{v}H_3)} \circ e^{\text{ad}(\frac{1}{v}H'_1)} \circ e^{\text{ad}(\frac{1}{v}H_2)} \\ &= \widetilde{\psi_1 \circ \psi_2} \circ e^{\text{ad}(\frac{1}{v}H(\psi_1, H_1, \psi_2, H_2))},\end{aligned}$$

where we used

(i) $\text{Diff}(M, \omega) \times \text{Aut}(C_M) \rightarrow \text{Aut}(C_M)$, smooth,

$$(\psi, \Psi) \longmapsto \tilde{\psi}^{-1} \circ \Psi \circ \tilde{\psi},$$

(ii) $\text{Diff}(M, \omega) \times \text{Diff}(M, \omega) \rightarrow \text{Aut}(C_M)$, smooth.

$$(\psi, \phi) \longmapsto (\widetilde{\psi \circ \phi})^{-1} \circ \tilde{\psi} \circ \tilde{\phi}.$$

Obviously, $\text{Aut}(C_M)$ is a Lie sub group of $\text{Aut}(C_M)$.

Using the following sequence

$$1 \rightarrow \underline{\text{Aut}}(C_M) \rightarrow \text{Aut}(C_M) \rightarrow \text{Diff}(M, \omega) \rightarrow 1$$

$\tilde{\phi} \longleftarrow \phi$

we complete the proof. ■